Lectures

Dwarf Galaxies & Dark-Dark Halos: Feedback Processes

- · The "fundamental line"
- Origin of scaling relations: supernova feedback
- Dark-dark halos (DDH) must exist
- Origin of DDH by photoionization
- Halo substructure: phase-space density

1. Missing Dwarfs & the "Fundamental Line"

Dekel & Woo 2003







z=49.000 Halo buildup by mergers



Λ CDM model: many dwarf satellites



Moore et al



Low-Surface-Brightness & Dwarf Galaxies

197

Only a few faint dwarf satellites



Antlia (Tr)





The "Fundamental Line" of LSB/Dwarfs



Metallicity



Local Group Dwarfs: Metallicity





2. Origin of Scaling Relations:

virial theorem & spherical halo collapse angular momentum feedback



Bright Galaxies: Tully Fisher Relation



Courteau et al. 04

Surface Brightness: SDSS



Metallicity SDSS



Supernova Feedback Dekel & Silk 86 Dekel & Woo 03



Much energy in SNe



Galactic wind M82





Chandra X-Ray Observatory image of M82

Simulation of supernova blowout





Mori et al

Supernova Feedback



Fragile, Murray, Lin 04

Supernova Feedback Scale

(Dekel & Silk 86)

Energy fed to the ISM during the "adiabatic" phase:

$$E_{\rm SN} \approx v \varepsilon \ \dot{M}_* \ t_{\rm rad} \propto M_* (t_{\rm rad}/t_{\rm ff})$$

$$\dot{M}_* \approx M_*/t_{\rm ff} \qquad \approx 0.01$$

for $\Lambda \propto T^{-1}$ at $T \sim 10^5 K$

Energy required for blowout:

$$E_{\rm SN} \approx M_{\rm gas} V^2$$

$$\rightarrow V_{\rm crit} \approx 100 \text{ km/s} \rightarrow M_{*\rm crit} \approx 3 \times 10^{10} M_{\odot}$$



LSB vs HSB





The "Fundamental Line" of LSB/Dwarfs



Metallicity



Local Group Dwarfs: Metallicity



LG Dwarfs: Velocity



Summary: SN feedback

Could be responsible for the transition scale at $M_*=3\times10^{10}$, and the "fundamental line" of LSB/dwarf galaxies, $M^*/M^{\infty}V^2$.

A lower bound for galaxies







It isn't that simple to turn on the light



3. Dark-Dark Halos Must Exist





Dark Dark Halos must exist!

Search for DDH

Complete removal of gas from proto-halos?

By SN outflow? unlikely

By ram pressure due to outflow from a nearby galaxy (Scannapieco, Ferrara & Broadhurst 00)? By radiative feedback?

4. Evaporation by Thermal Winds Shaviv & Dekel 2003

Radiative Feedback

Reionization of H by UV flux from stars and AGN by $z_{ion} \sim 10 \rightarrow$ heating gas to $T \approx (1-2) \times 10^4$ K Jeans scale – no infall into halos of V<30 km/s Efstathiou 92; Thoul & Weinberg 96; Gnedin & Ostriker 97; Gnedin 00

But complete gas removal? Evaporation from halos of V<10 km/s Barkana & Loeb 99 V<30 km/s Shaviv & Dekel 03

May eliminate luminous dwarfs in small halos, 10<V<30

Evaporation of hot gas

cold gas hot gas

Mass loss from top of potential well $t_{evap} \approx t_{dyn} e^{\phi/kT}$ It is continuously replenished and lost Continuous energy input by the ionizing flux \rightarrow steady wind

Steady Thermal Wind In stars: Parker 1960. In galaxies: extended potential well

Hydrodynamics:

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{v})$$
$$\rho \frac{D \mathbf{v}}{D t} = -\nabla P + \mathbf{f}_{\text{grav}}$$
$$P = c_{\text{s}}^{2} \rho$$

Assume: spherical, c_s=const., steady state $\frac{\dot{M}(r) = \text{const.} \rightarrow \dot{\rho} = 0}{\dot{N}(r) = \text{const.} \rightarrow \dot{\rho} = 0}$

→ wind equation: $\left(v(r) - \frac{c_s^2}{v(r)} \right) v'(r) = -\phi'(r) + \frac{2c_s^2}{r}$

→ the sonic radius: $\phi'(r_s) = 2c_s^2 / r_s$ → $r_s \approx GM / c_s^2$

wind parameter (NFW): $\psi = \frac{GM_{c}/r_{c}}{c_{c}^{2}}$

$$t_{evap} / t_{dyn} \approx 10^{\psi - 1}$$

 ψ >>1 tightly bound, no evaporation ψ >1 bound, but steady wind >> t_{dyn} ψ <1 rapid evaporation ~ t_{dyn}

Summary Dwarf Halos Dark-dark halos must exist at V<30 km/s Half the photo-ionized gas evaporates by steady winds from halos of V<30 km/s. Halos in the range 10<V<30 could be: - gas-poor dSph /dE - or totally dark

No galaxies V<10 becasue of cooling barrier

Summary: Characteristic Scales

	V (km/s)	$M_{\star}(M_{})$	M(M_)	<u> </u>
Cooling (Brems.)	300	2×10 ¹¹	10 ¹³	clusters
Shock heating	100	3×10 ¹⁰	6×10 ¹¹	L*, Tyoung disks
Supernovae	100	3×10 ¹⁰	6×1011	LSB T
Photoionization	30	10 ⁸	2×10 ¹⁰	d <u>Sp</u> h d <u>ar</u> k
Cooling (H)	10	3×10 ⁵	6×10 ⁸	

Phase-Space Density & Halo Substructure

Arad & Dekel, in progress

Distribution function of f:

$$V(f = f_0) \equiv \int d\vec{x} d\vec{v} \,\delta_{Dirac}[f(\vec{x}, \vec{v}, t) - f_0]$$

V(f)df = volume of phase space occupied by f in the range (f,f+df)

Measuring f(x,v) using an adaptive "grid" Delaunay Tesselation

$$f_i = (d+1)\frac{m}{V_i}$$

Arad, Dekel & Klypin

PDF of Phase-Space Density

Arad, Dekel & Klypin

PDF of Phase-Space Density

f

Arad, Dekel & Klypin

V(f) related to _(r)?

e.g., spherical & isotropic

$$\rho(r) \propto r^{-\alpha}, \quad V(f) \propto f^{-\beta}, \quad \beta = \frac{18 - 4\alpha}{6 - \alpha}$$

$$\alpha = 3 \Leftrightarrow \beta = 2$$

$$\alpha = 2 \iff \beta = 2.5$$

$$\alpha = 1 \iff \beta = 2.8$$

$$\alpha = 0 \iff \beta = 3$$

Halo Phase-Space Density

Real Density

Phase-Space Density

Halo Phase-Space Density

Real Density

Phase-Space Density

Profiles in Real Space and Phase Space f(r)

radius

Is $v(f) \propto f^{-2.5}$ determined by substructure?

ACDM No short waves

Real-Space Density

Moore et al.

Phase-Space density

$\land CD$

No short waves

Phase-Space Density Profile

 $\land CD$

No short waves

Same power law v(f)?

Additive Contribution of Subhalos

The Two Most Massive Subhalos

Adding up Sub-halos

Tentative Conclusions

In hierarchical clustering, robust PDF: v(f)∝f^{-2.5} doesn't depend on power-spectrum slope, or on method of simulation The power-law v(f) is driven by substructure. How exactly? Yet to be understood !

Phase-space density is a unique tool for studying substructure and its evolution

Adding up small CDM halos leads to $v(f) \propto f^{-2.5}$? How robust? How dependent on subhalo density profile and mass function?