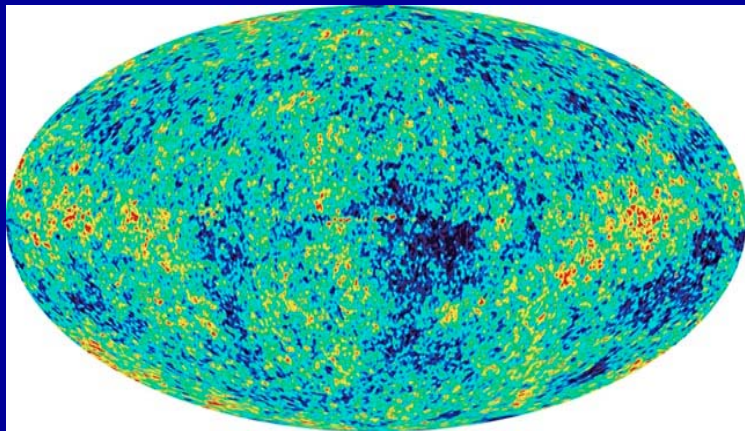
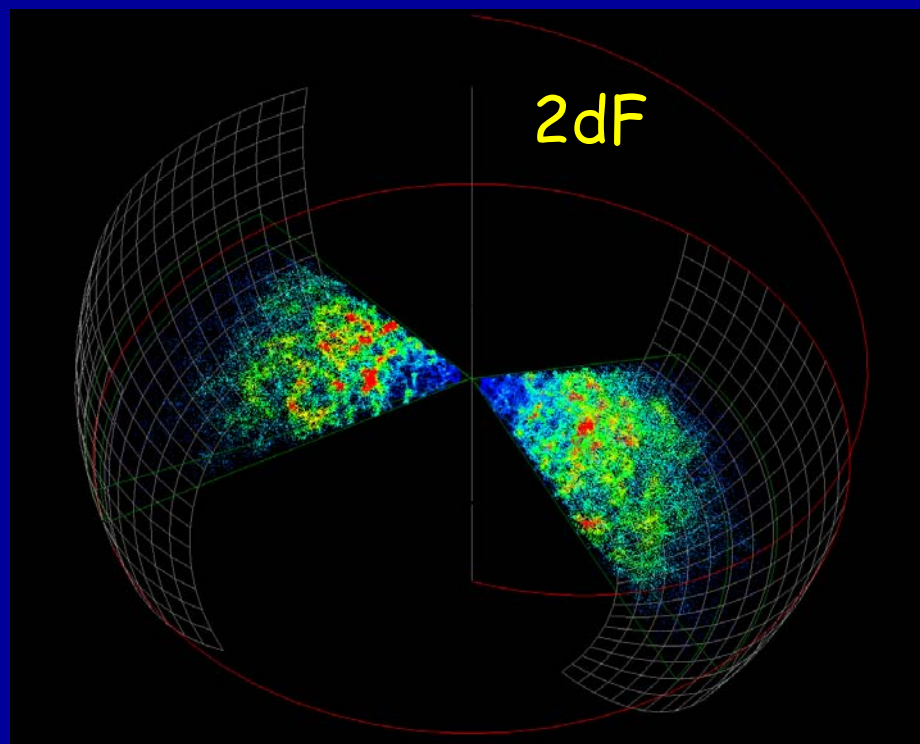


Linear Growth of Fluctuations by Gravitational Instability

WMAP



2dF

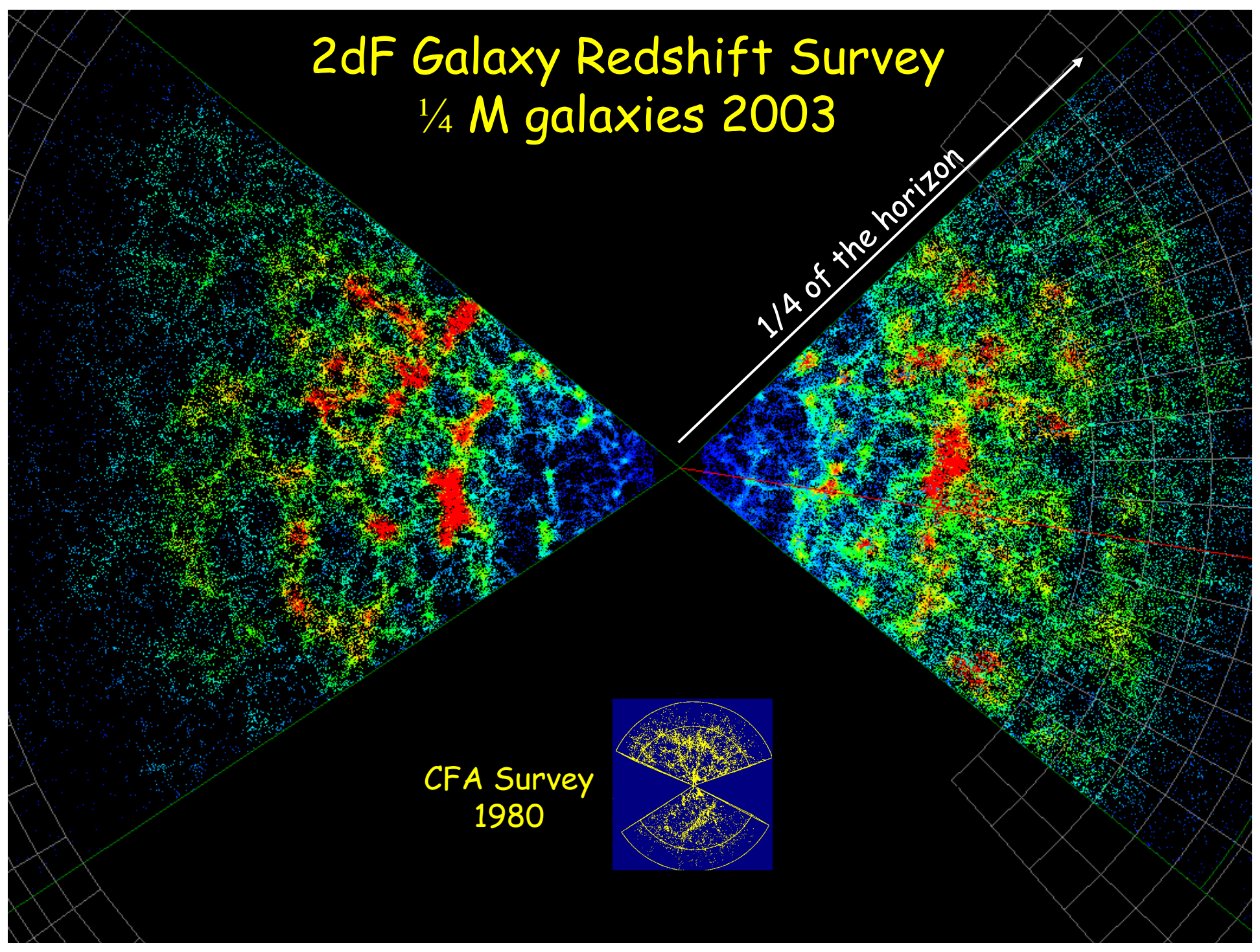
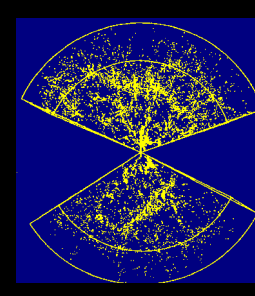


2dF Galaxy Redshift Survey

$\frac{1}{4}$ M galaxies 2003

$\frac{1}{4}$ of the horizon

CFA Survey
1980





DAWN
OF
TIME

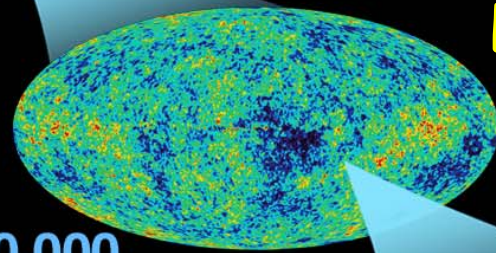
H₀

Early Universe

tiny fraction
of a second



inflation

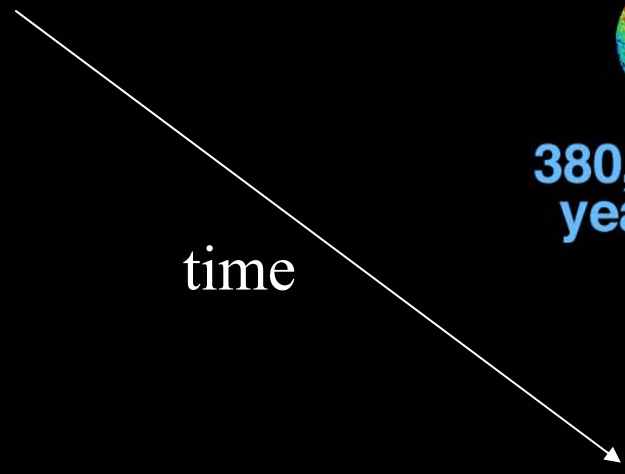


Cosmic Microwave
Background

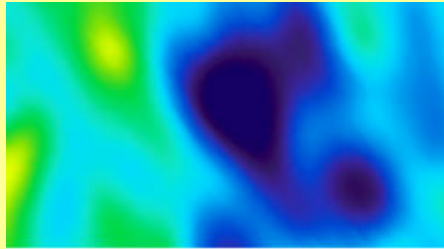
380,000
years

Today

13.7
billion
years



Late Cosmological Epochs



380 kyr $z \sim 1000$

recombination
last scattering

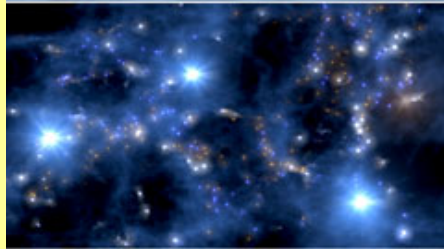


dark ages

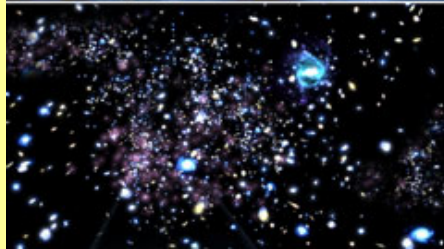


180 Myr $z \sim 20$

first stars
reionization



galaxy formation

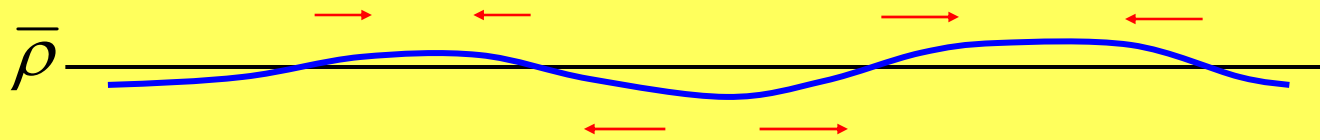


13.7 Gyr $z=0$

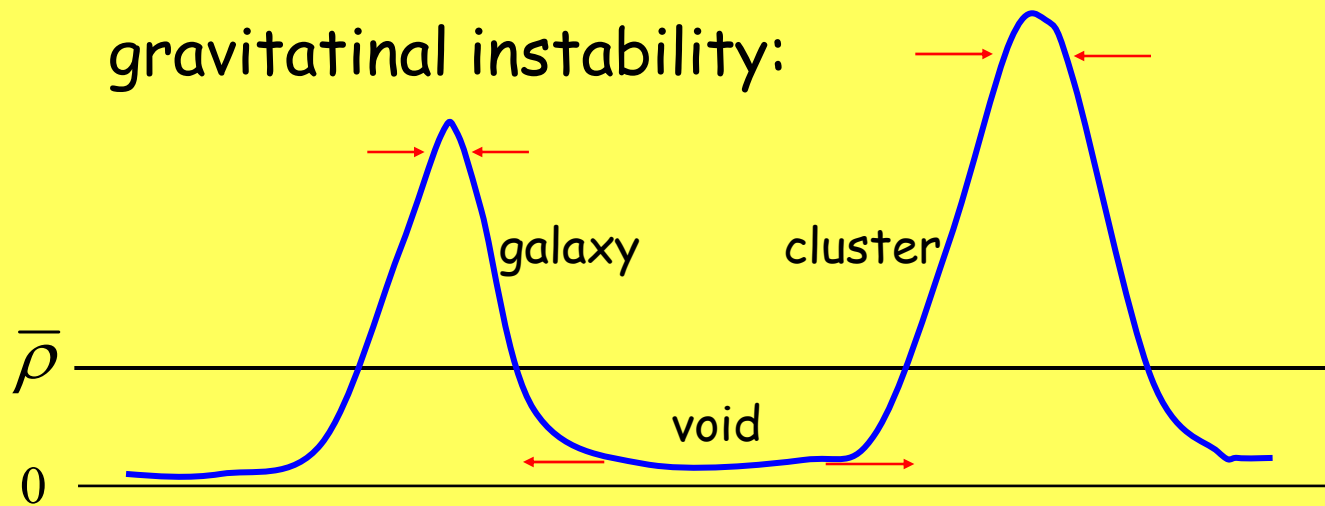
today

Gravitational instability

small-amplitude fluctuations:



gravitational instability:



Gravitational Instability: linear, matter-era

Fluid equations :

$$(1) \quad \dot{\rho} + \vec{\nabla} \cdot (\rho \vec{V}) = 0 \quad \text{continuity}$$

$$(2) \quad \dot{\vec{V}} + (\vec{V} \cdot \vec{\nabla}) \vec{V} = -\vec{\nabla} \Phi - \vec{\nabla} P / \rho \quad \text{Euler}$$

$$(3) \quad \nabla^2 \Phi = 4\pi G \rho \quad \text{Poisson}$$

Uniform background : $\rho(\vec{r}) = \text{const.}$

$$\vec{r} \equiv a\vec{x} \quad \vec{v} = \frac{\dot{a}}{a} \vec{r} \quad \rho^{(1)} = \frac{\rho_0}{a^3} \quad \frac{\dot{a}^2}{a^2} = \frac{8\pi G}{3} \rho \quad \frac{\ddot{a}}{a} = -\frac{4\pi G}{3} \rho$$

$$H \equiv \dot{a} / a \quad \dot{\rho}^{(1)} = -3\rho H$$

Perturbations : $\rho(\vec{r}, t) = \rho_u(t)[1 + \delta(\vec{r}, t)] \quad \vec{V} = H(t)\vec{r} + \vec{v} \quad \Phi = \Phi_u + \varphi \quad P = p$

1st order + : $\delta \ll 1$ etc.

$$(1) \quad \dot{\delta} + H\vec{r} \cdot \vec{\nabla} \delta + \vec{\nabla} \cdot \vec{v} + \vec{\nabla} \cdot (\delta \vec{v}) = 0$$

$$(2) \quad \dot{\vec{v}} + H(\vec{r} \cdot \vec{\nabla}) \vec{v} + H\vec{v} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = -\vec{\nabla} \varphi - c_s^2 \vec{\nabla} \delta$$

$$(3) \quad \nabla^2 \varphi = 4\pi G \rho_u \delta$$

$$\vec{\nabla} P = \frac{\partial P}{\partial \rho} \vec{\nabla} \rho \quad c_s^2 \equiv \frac{\partial P}{\partial \rho} (\text{ideal gas}) = \frac{P}{\rho} = \frac{kT}{m_p}$$

Comoving coordinates: $\vec{x} \equiv \frac{\vec{r}}{a}$ $\frac{\partial}{\partial t}\Big|_x = \frac{\partial}{\partial t}\Big|_r + \frac{\dot{a}}{a} \vec{r} \cdot \vec{\nabla}_r \Big|_t$ $\nabla_x = a \nabla_r$

$\vec{w} \equiv \vec{v} / a$ $\psi \equiv \varphi / a$

(1) $\dot{\delta} + \vec{\nabla} \cdot \vec{w} + \vec{\nabla} \cdot (\delta \vec{w}) = 0$

(2) $\dot{\vec{w}} + 2H\vec{w} + (\vec{w} \cdot \vec{\nabla})\vec{w} = -\vec{\nabla} \psi - a^{-1} c_s^2 \vec{\nabla} \delta$

(3) $\nabla^2 \psi = 4\pi G \rho_u \delta$ [= (3/2)H²Ωδ]

$a^{-1} \vec{\nabla} \cdot (\text{eq. 2})$ $\partial / \partial t (\text{eq. 1}) \rightarrow$

$\ddot{\delta} + 2H\dot{\delta} = 4\pi G \rho_u \delta + a^{-2} c_s^2 \nabla^2 \delta$

gravity pressure

$\delta(\vec{x}, t)$ $a(t)$ $H(t) = \frac{\dot{a}}{a}$ $\rho_u(t) \propto a^{-3}$

$$\ddot{\delta} + 2H\dot{\delta} = 4\pi G\rho_u \delta + a^{-2}c_s^2 \nabla^2 \delta$$

$$a(t) \quad H(t) = \frac{\dot{a}}{a} \quad \rho_u(t) \propto a^{-3}$$

Static background : $\dot{a} = 0 \quad a = \text{const.} \equiv 1 \quad \rho_u = \text{const.}$

$$\delta \propto \exp[i(\vec{k} \cdot \vec{x} + \omega t)] \quad \omega^2 = k^2 c_s^2 - 4\pi G\rho_u$$

pressure gravity

Jeans scale: $k_J = \left(\frac{4\pi G\rho_u}{c_s^2} \right)^{1/2} \quad \lambda_J \equiv \frac{2\pi}{k_J} \quad M_J \equiv \frac{4\pi}{3} \rho_m \left(\frac{\pi c_s^2}{G\rho} \right)^{3/2} \propto \frac{T^{3/2}}{\rho^{1/2}}$

$$\lambda \gg \lambda_J \quad (p=0) \quad \rightarrow \delta = Ae^{\omega t} + Be^{-\omega t}$$

$$\lambda \ll \lambda_J \quad \rightarrow \text{stable oscillations}$$

Expanding background , $\lambda \gg \lambda_J$:

$$k = 0 \quad \rightarrow \quad a \propto t^{2/3} \quad \rightarrow \quad \ddot{\delta} + \frac{4}{3t}\dot{\delta} = \frac{2}{3t^2}\delta \quad \rightarrow \quad \delta = At^{2/3} + Bt^{-1}$$

$$k = -1 \quad \rightarrow \quad a \propto t \quad \rightarrow \quad \ddot{\delta} + \frac{2}{t}\dot{\delta} = \frac{3\Omega_0 t_0}{2t^3}\delta \quad \rightarrow \quad \delta = \text{const.} \quad \text{freezout}$$

Properties of the linear growing mode:

linear $\ddot{\delta} + 2H\dot{\delta} = (3/2)H^2\Omega\delta$ $H(t)$ $\Omega(t)$

growing mode: $\delta \propto D(t)$

$$f(\Omega) \equiv \frac{\dot{D}}{HD} \approx \Omega^{0.6} \quad \rightarrow \quad \frac{\ddot{D}}{D} = H^2 \left(-2f + \frac{3}{2}\Omega \right)$$

continuity $\rightarrow \delta = -\frac{1}{Hf} \vec{\nabla} \cdot \vec{v}$

Poisson $\rightarrow \vec{v} = -\vec{\nabla} \varphi_v$ irrotational $\varphi = \frac{3H\Omega}{2f} \varphi_v$

The Jeans scale in an expanding universe:

In k -space: $\delta = \sum_{\vec{k}} \delta_{\vec{k}}(t) e^{i\vec{k} \cdot \vec{x}}$ $r = ax$

for each \vec{k} : $\ddot{\delta}_{\vec{k}} + 2H\dot{\delta}_{\vec{k}} = (4\pi G\rho - k^2 c_s^2) \delta_{\vec{k}}$ \rightarrow same Jeans scale

Lecture 3

Statistics of Fluctuations: The Cold Dark Matter Scenario

The Initial Fluctuations

At Inflation: Gaussian, adiabatic

fluctuation field $\delta(x) = \frac{\rho(x) - \langle \rho \rangle}{\langle \rho \rangle}$ a realization of an ensemble
ensemble average \sim volume average

Fourier $\delta(\vec{x}) = \sum_{\vec{k}} \delta_{\vec{k}} e^{i\vec{k}\cdot\vec{x}}$ Power Spectrum $P(k) \equiv \langle |\tilde{\delta}(\vec{k})|^2 \rangle \propto k^n$

rms

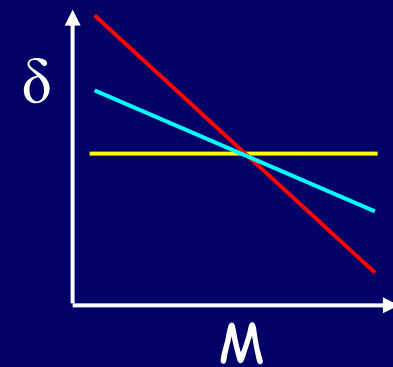
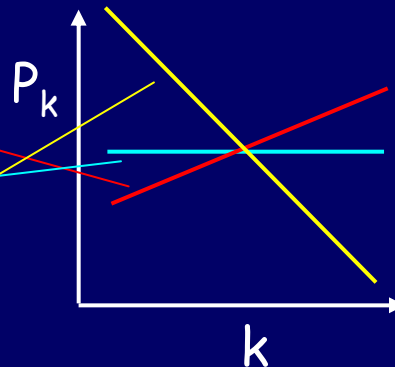
$$\langle \delta^2 \rangle_\lambda \sim \left\langle \int_{k=0}^K \int_{k'=0}^{K-2\pi/\lambda} \exp[-i(k+k')\cdot x] d^3k' d^3k \delta_k \delta_{k'} \right\rangle \sim \int_{k=0}^K d^3k \langle \delta_k \delta_{-k} \rangle$$

$\longleftarrow \delta_{Dirac}(k+k') \longrightarrow$

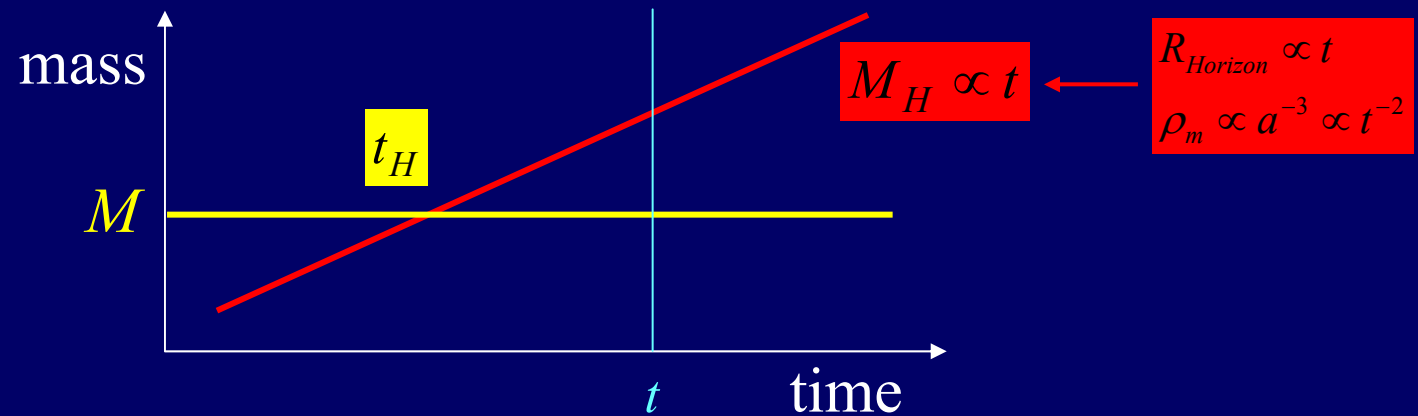
$$\langle \delta_k \delta_{-k} \rangle = \langle |\delta_k|^2 \rangle$$

$$\langle \delta^2 \rangle_\lambda \propto \int_{k=0}^{2\pi/\lambda} P_k d^3k \propto M^{-(n+3)/3}$$

$n = 1$	$\delta \propto M^{-2/3}$
$n = 0$	$\delta \propto M^{-1/2}$
$n = -3$	$\delta \propto const.$

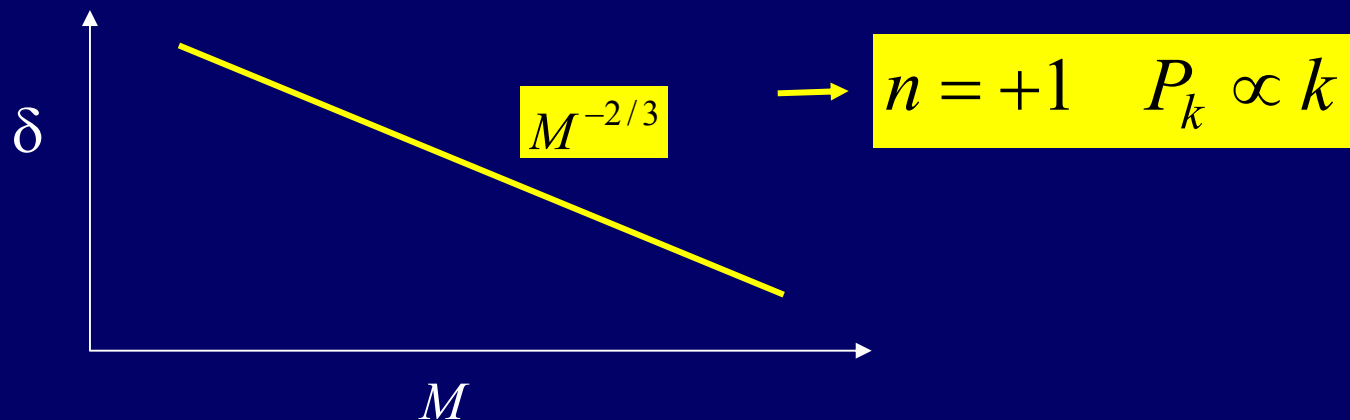


Scale-Invariant Spectrum (Harrison-Zel'dovich)



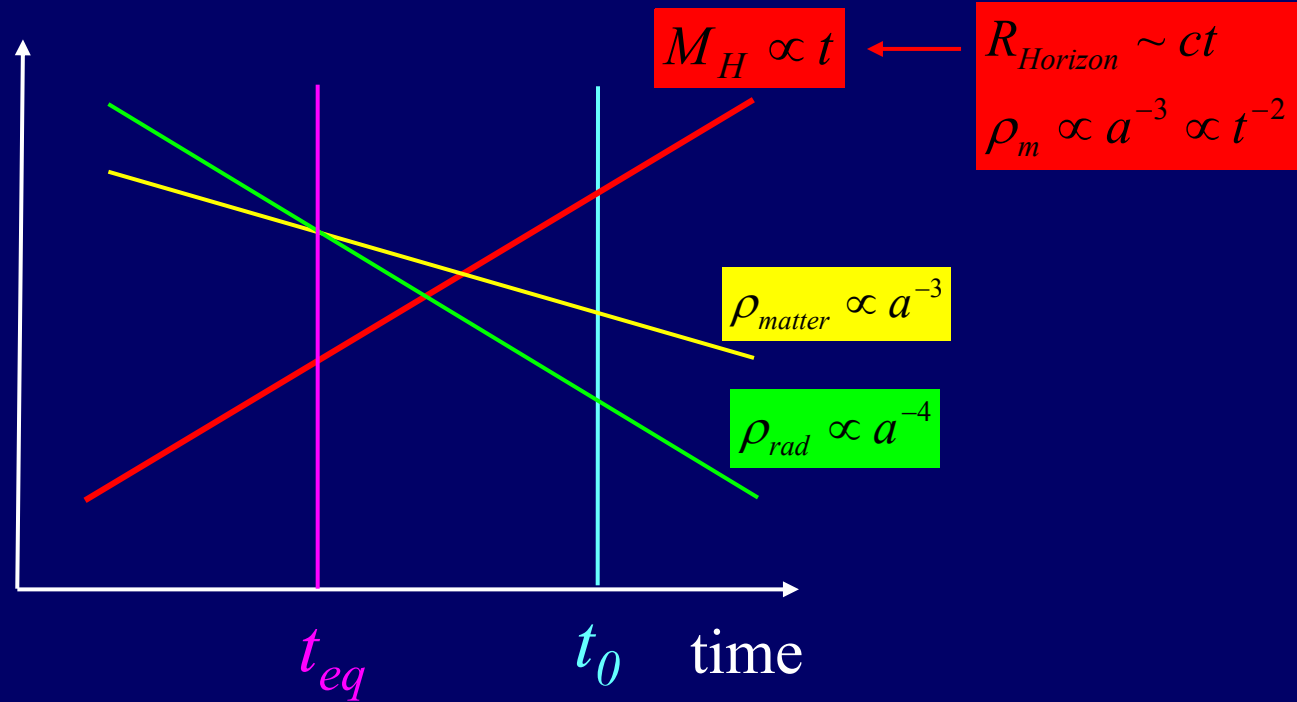
$$\delta(M, t) = \delta_H \left(\frac{t}{t_H(M)} \right)^{2/3} \propto M^{-2/3} t^{2/3}$$

$\delta_H = \text{const.}$



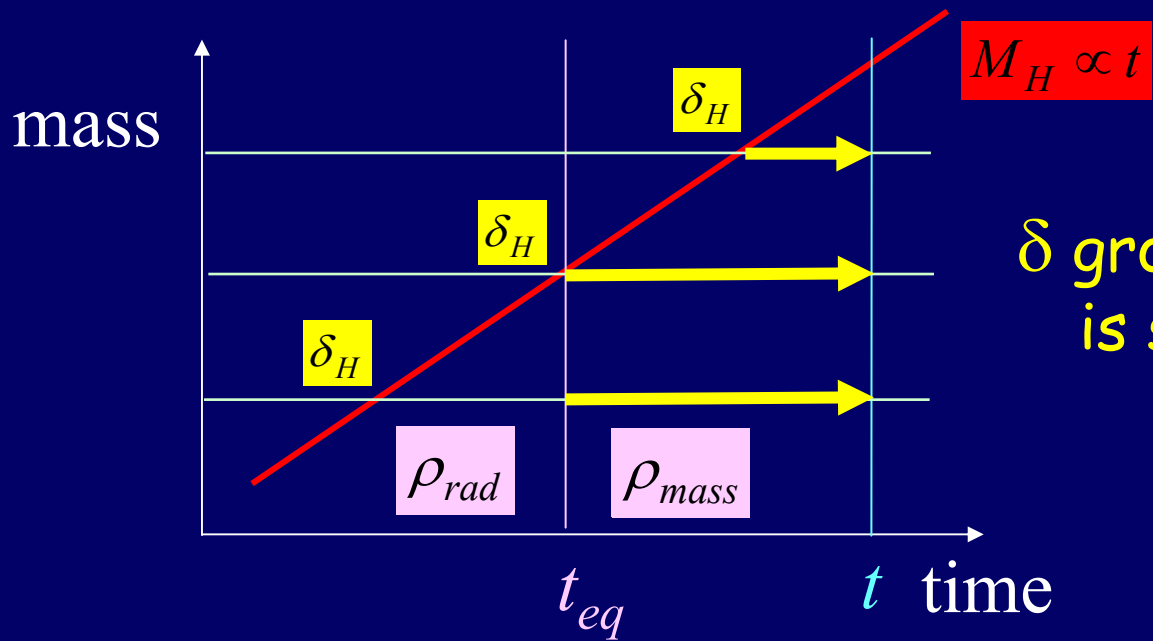
Cosmological Scales

mass

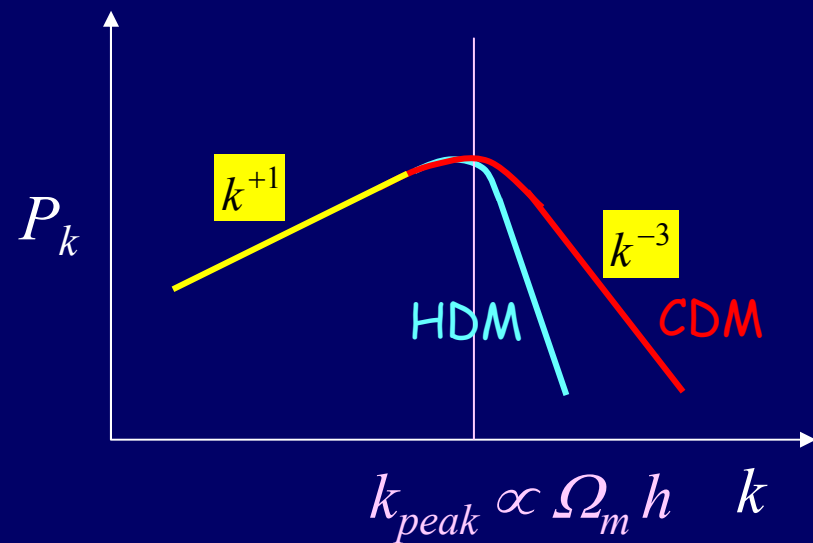
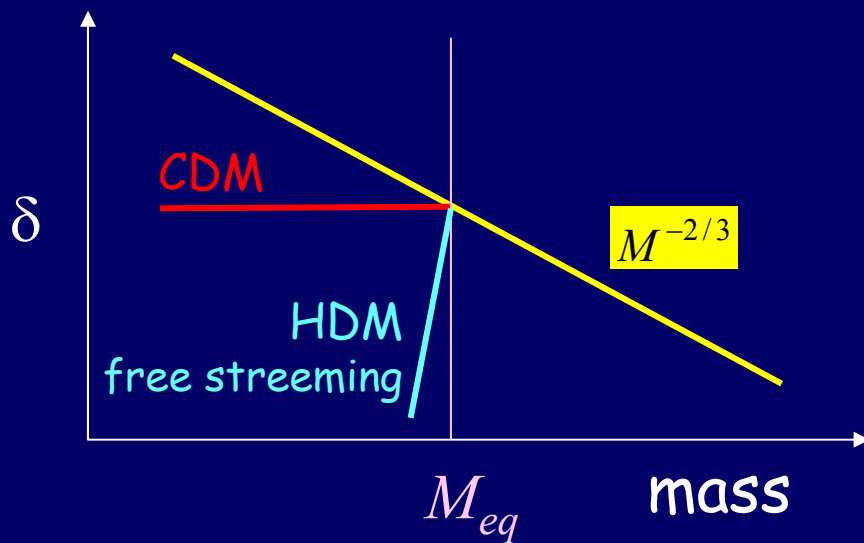


$$z_{eq} \sim 10^4$$

CDM Power Spectrum



δ growth when matter is self-gravitating



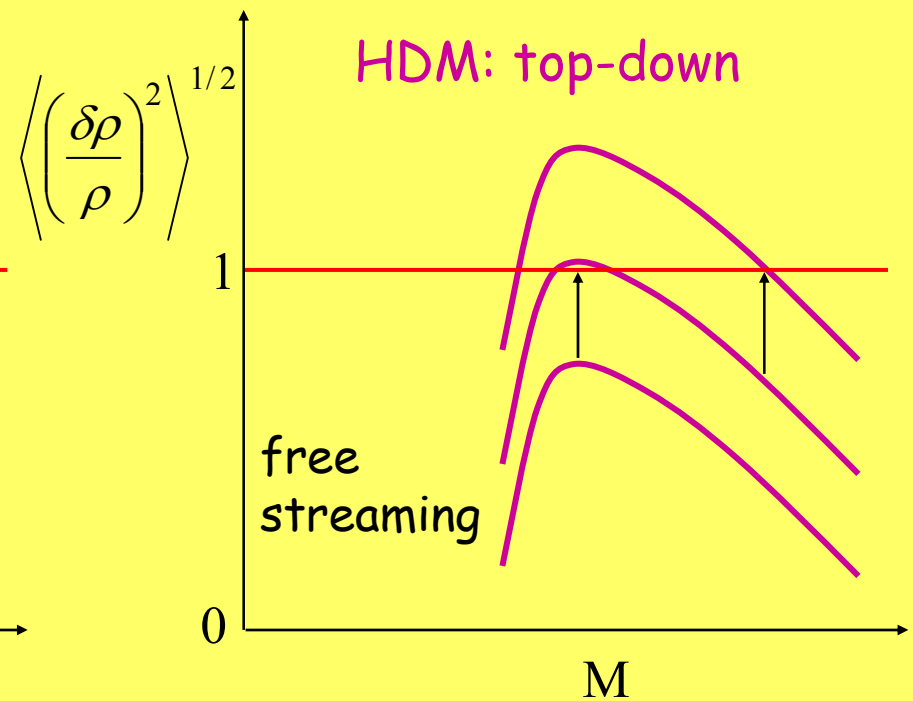
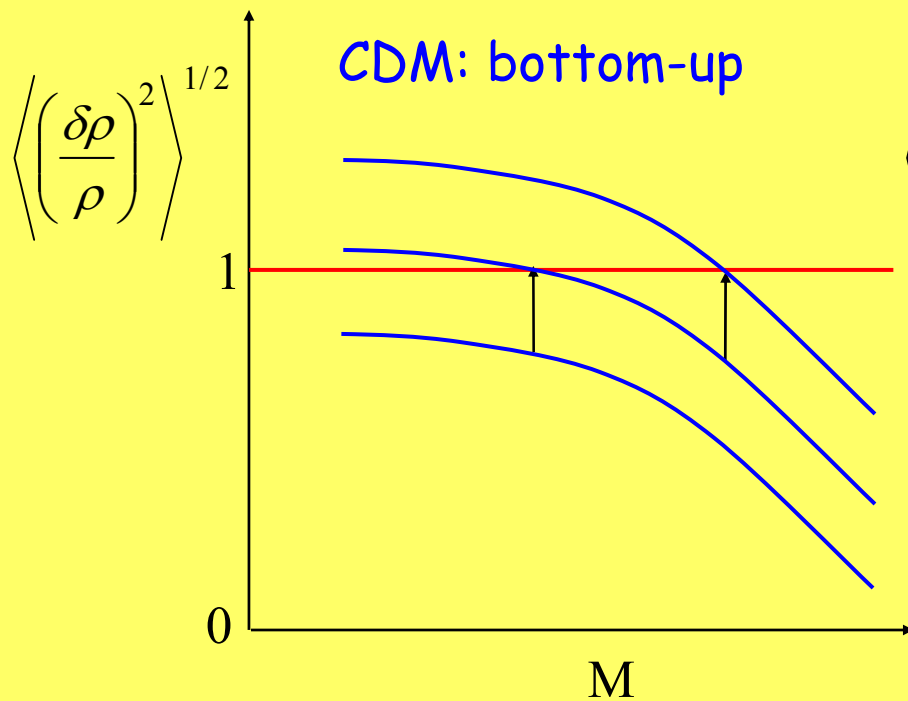
Formation of Large-Scale Structure

Fluctuation growth in the linear regime: $\delta \ll 1 \rightarrow \delta \propto a \propto t^{2/3}$

rms fluctuation at mass scale M : $\delta \propto M^{-\alpha} \quad 0 < \alpha = (n+3)/6 \leq 2/3$

Typical objects forming at t : $1 \sim \delta \propto M^{-\alpha} a \rightarrow M_* \propto a^{1/\alpha}$

example $n = -2 \rightarrow M_* \propto a^6$



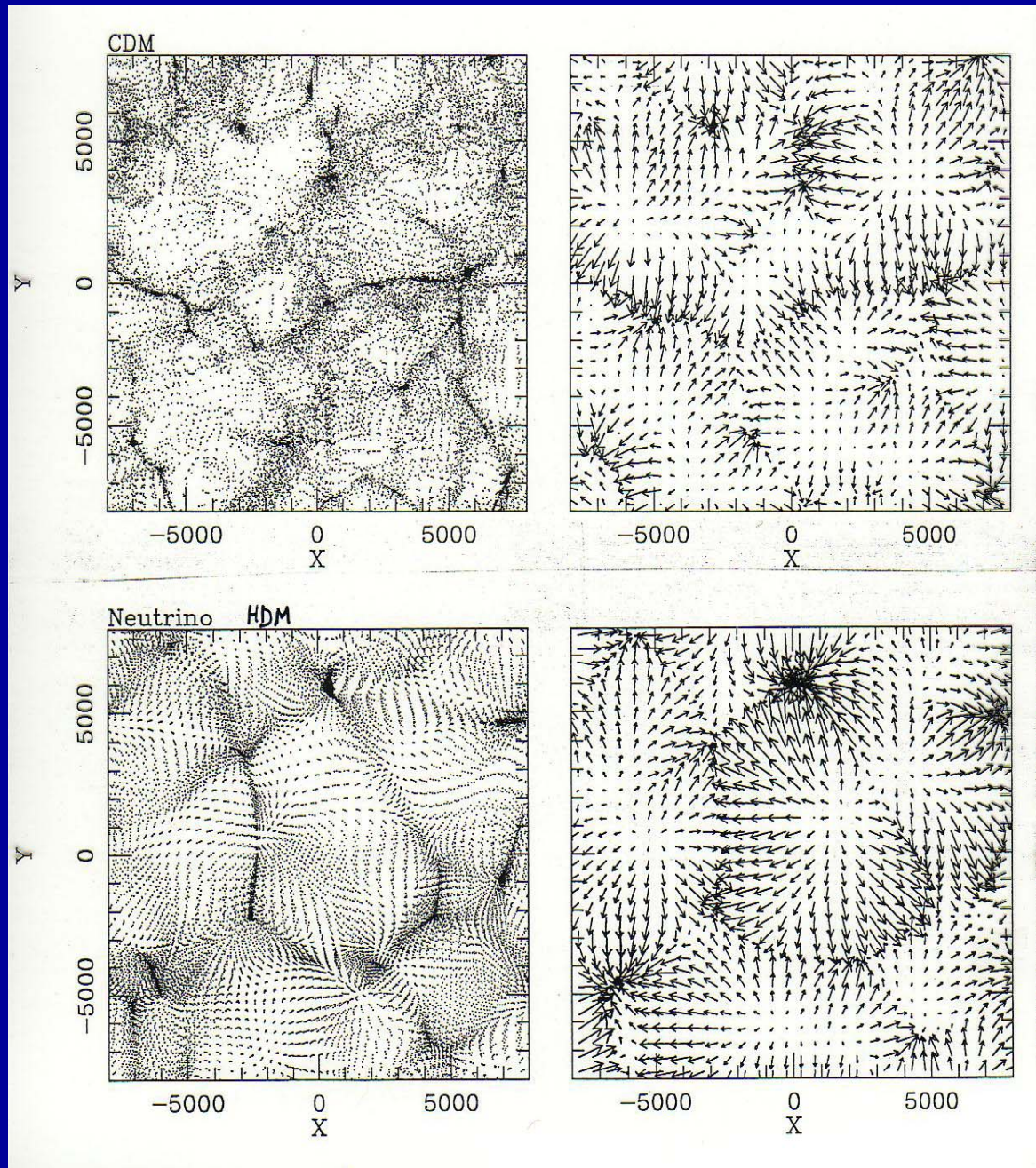
Micro-Macro Connection

Cold Dark Matter

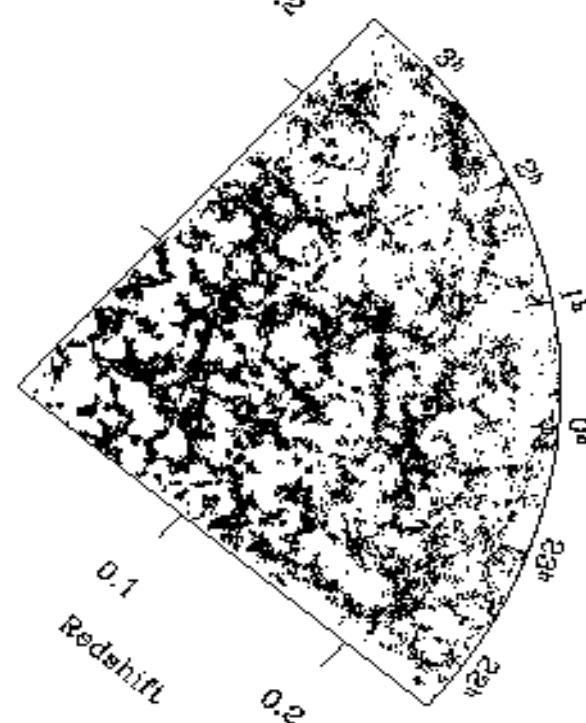
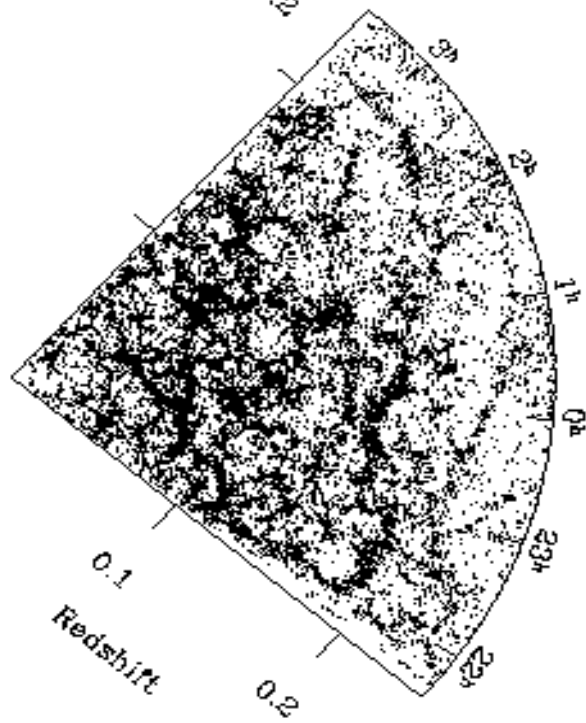
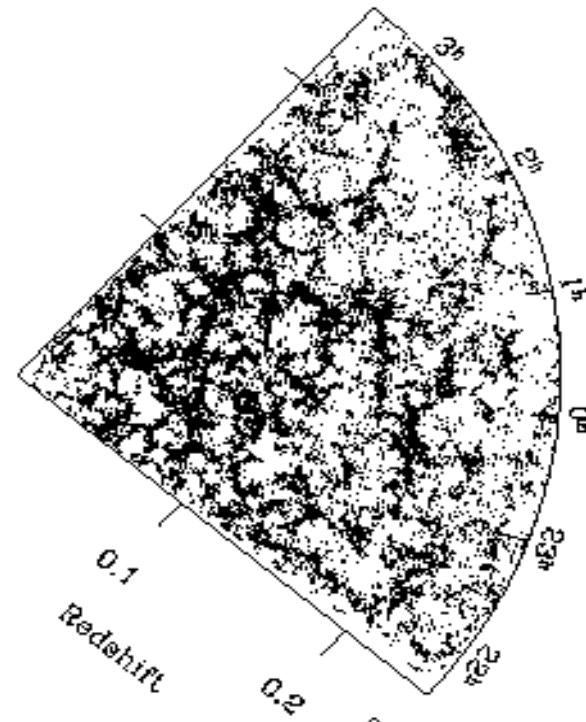
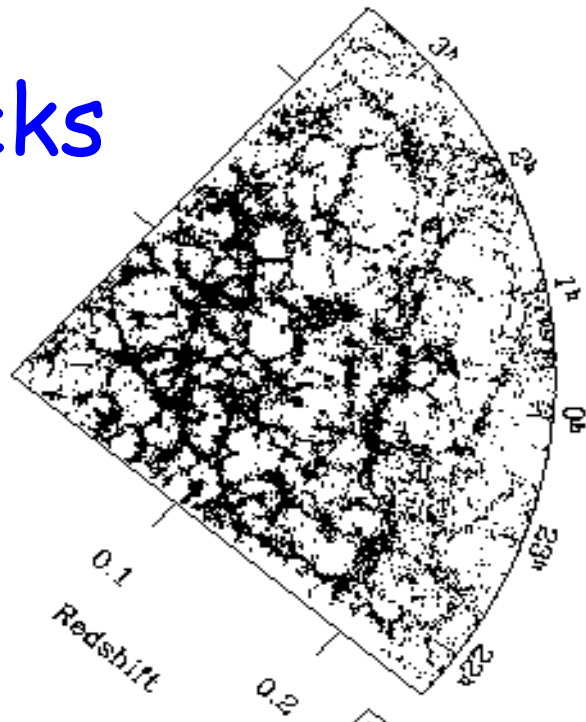
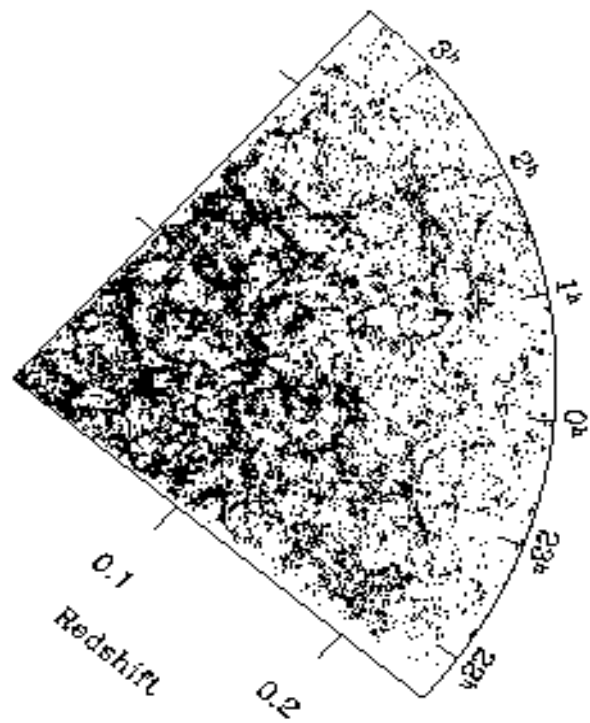


Hot Dark Matter

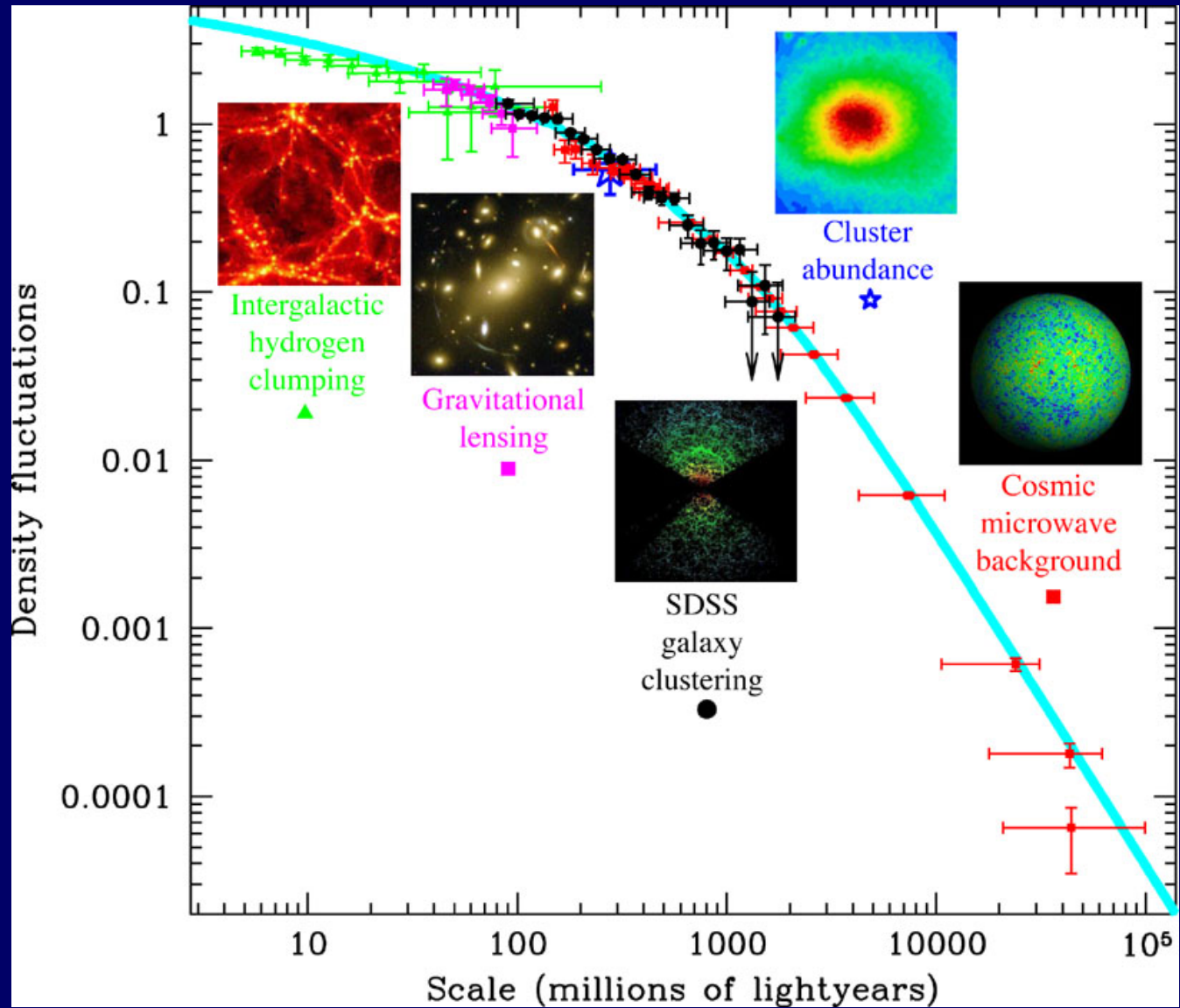
ν



2dF and Mocks



Power Spectrum



Λ CDM Power Spectrum

$$P(k) \propto k T^2(k)$$

$$T(k) = \frac{\ln(1 + 2.34q)}{2.34q} \left(1 + 3.89q + (16.1q)^2 + (5.46q)^3 + (6.71q)^4\right)^{-1/4} \quad q = \frac{k}{\Omega_m h^2 \text{Mpc}^{-1}}$$

normalization: $\sigma_8 \equiv \sigma_{\text{tophat}}(R = 8h^{-1} \text{Mpc})$