

Interactions amongst particles:
Emergence
in
condensed matter
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<http://physics.ucsc.edu/~sriram/papers/>

Quantum theory of many particles:

Interactions add spice

- Bosons and Fermions
- Atomic physics- beyond the Hydrogen atom- Hartree Fock theory- scaling-ideas
- Bosons: superfluidity- beyond the Bose condensation of ideal gas - probed by atomic traps and cold atom experiments.
- Solids:
 - ★ Bands and Hartree Fock theory
 - ★ Landau Fermi Liquid theory
 - ★ Magnetism
 - ★ Superconductivity
 - ★ Mott Physics
 - ★ Extreme correlations
- *New models- exactly solvable type

Main thrust is on emergent properties of many particle systems.

Emergence = new and fundamental laws arising from assemblies of particles that are unexpected from single particle properties- e.g. superconductivity, superfluidity etc.

Emergence plays a big role in Condensed Matter Physics and also in Quantum field theory.

Wavefunctions for N particles

$$\psi(x_1, \dots, x_i, \dots, x_j, \dots, x_N)$$

$$\psi(x_1, \dots, x_i, \dots, x_j, \dots, x_N) = \pm \psi(x_1, \dots, x_j, \dots, x_i, \dots, x_N) \quad \begin{array}{l} \text{Bose} \\ \text{Fermi} \end{array}$$

Bosons in Nature:

- Elementary particles: pions,...all integer spin particles
- More useful idea: even number of Fermions act as bosons. Fermions are e,n,p,....
- ^4He (2 p+ 2n)-Kammerling Onnes 1908 liquefied and in 1938 Kapitza found superfluid property
- Atoms Rubidium-87, Sodium-23 .. Cooling by atomic traps and condensation (Nobel Prize 1997, 2001)

From Atoms to Solids: Hartree Fock theory and Landau Fermi Liquids

$$H = -\frac{\hbar^2}{2m} \sum_i \nabla_i^2 + \sum_{i < j} \frac{e^2}{|r_i - r_j|} + \sum_i U(r_i)$$

U(r) is say Hydrogenic atomic potential

Let the mean separation between electrons be called r_s

$$\frac{4\pi}{3} r_s^3 = \frac{V}{N} = \frac{1}{n}$$

We next scale all lengths by r_s

$$\nabla^2 \rightarrow \frac{1}{r_s^2} \hat{\nabla}^2$$

$$r_i \rightarrow r_s \hat{r}_i$$

$$H = \frac{1}{r_s^2} \left(\hat{T} + r_s \hat{V} \right) + U$$

Therefore as r_s becomes small we can use perturbation theory around kinetic energy. This is at the heart of the Hartree Fock theory.

Question: What about Dirac eqn? (Graphene)

- Implications is that we can think of high density electrons as almost free electrons and treat potential energy as a perturbation.
- This is the reason why we can confidently predict the properties of heavy atoms with $Z > 3$, and also solids. For example in Aluminum, we have 10^{23} electrons per cc, excellent approximation is to use free electron theory- works like a charm,
- With Dirac type spectrum it is challenging though

Landau theory of Fermi liquids is a formalization of this idea.

In stat mech we learn that gas-> liquid are the same, liquids are more dense and hence more interacting. Similarly

Fermi gases <-> Fermi Liquids.... same and yet slightly different.

Fermi Liquids have certain standard set of properties, with “renormalization”

$$H_{eff} = \sum_{k\sigma} \varepsilon_k \delta n_{k\sigma} + \frac{1}{N} \sum_{kk'} f(k\sigma, k'\sigma') \delta n_{k\sigma} \delta n_{k'\sigma'}$$

Set of properties are controlled by Landau theory:

- Specific Heat $C_v = \gamma T + O(T^3)$
- Susceptibility $\chi = \text{constant}$
- Compressibility κ
- Quasi particle lifetimes $1/\tau \sim (T^2 + \omega^2)$
- Resistivity $\rho = \text{const} + T^2$

Large number of systems Liquid ^3He , many metals satisfy the diktats of Landau Fermi Liquid theory.

Probed by many means, including calorimetry, magnetic measurements, ARPES, transport.

These are the “Normal” systems

Now for the exceptions:

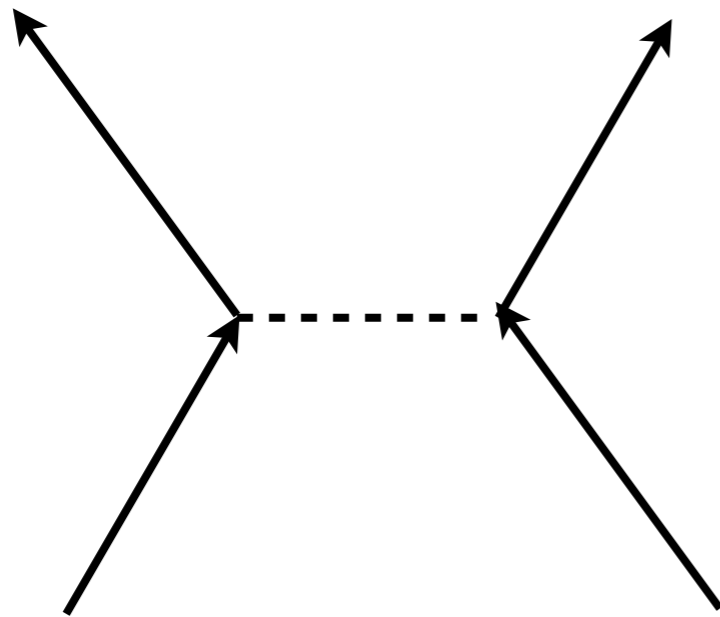
Liquids become solids at low temperature due to “broken symmetry” + interactions (Translation invariance in this case)

Fermi liquids become either magnets (of various denomination), or superconductors by breaking spin-rotation symmetry or gauge symmetry at low enough temperatures.

Superconductivity

$$H = \text{Kinetic Energy} - |U_{eff}| \sum_i n_{i\uparrow} n_{i\downarrow}$$

$$U_{eff} = U_{electron-phonon} - U_{Coulomb}$$



$$H = KE - |U_{eff}| \sum_{k,p} b_k^\dagger b_p;$$

$$b_k = C_{k\downarrow} C_{-k\uparrow}$$

$$\sum_k \langle b_k \rangle = \Delta$$

Violates gauge symmetry.

BCS theory of superconductors follows.

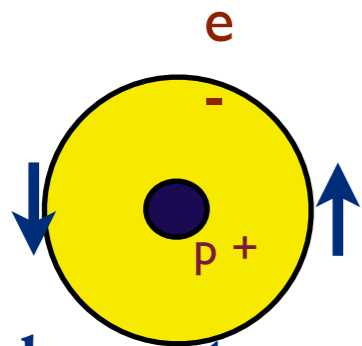
Attraction does not just create simple molecules- it can do a lot more. Macroscopic QM
 Non perturbative in U! One of the most comprehensive theories in Physics with detailed
 predictions for many measurements.

Mott Physics and Extreme Correlations

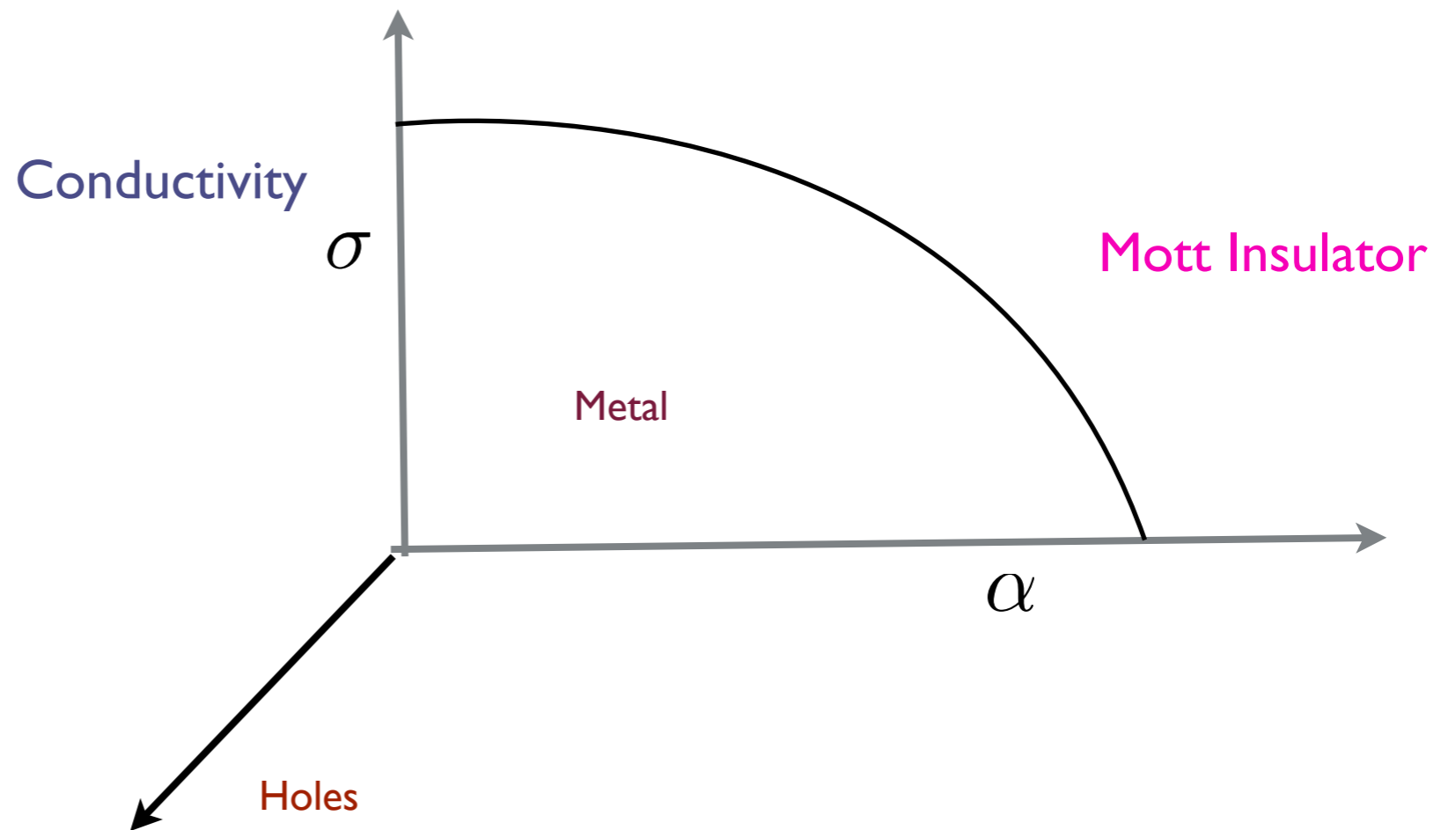
$$H = \text{Kinetic Energy} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

When U becomes larger than the range of the kinetic energy, we get a new possibility that was described by Mott. At exactly 1 electron per atom, we get an insulator at large U .

Array of Hydrogen Atoms
with separation α



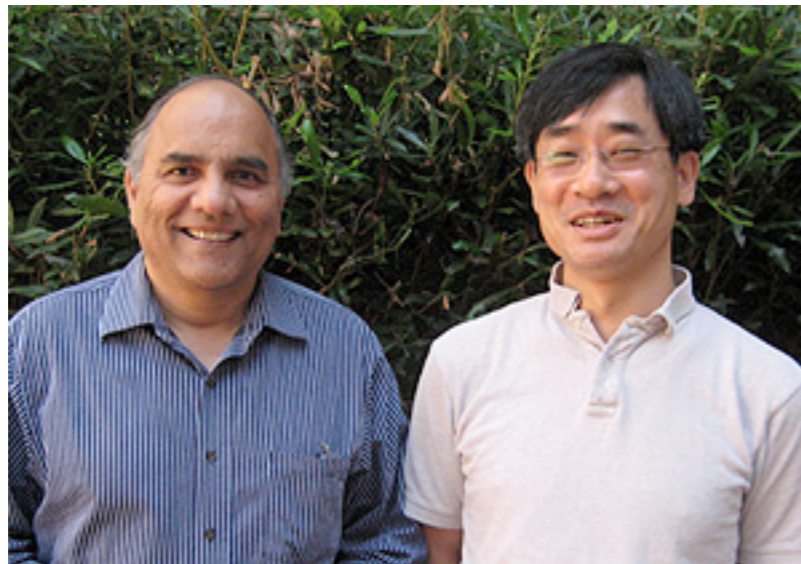
Hydrogen atom
The 1-s level may contain
0, 1 or 2 electrons



Description of the doped Mott Insulator is expected to lead to an understanding of High Tc systems- the holy grail!

The Physics of Extreme Correlations in Cuprates and its implications for ARPES

How I Learned to Stop Worrying and Love the Limit of Infinite U



Gey Hong

[1] Extremely Correlated Fermi Liquids, Phys. Rev. Letts. 107, 056403 (2011).

[2] Extremely Correlated Fermi Liquid Description of Normal State ARPES in Cuprates, Phys. Rev. Letts. 107, 056404 (2011). (G.-H. Gweon, B. S. Shastry and G. D. Gu.)

[3] "Anatomy of the self energy", Phys. Rev. B (2011)

[4] "Dynamical Particle Hole Asymmetry in Cuprate Superconductors", arXiv:1110.1032

ARPES = angle resolved photo emission

Significant probe since it measures the single electron propagator using the photoemission theory.

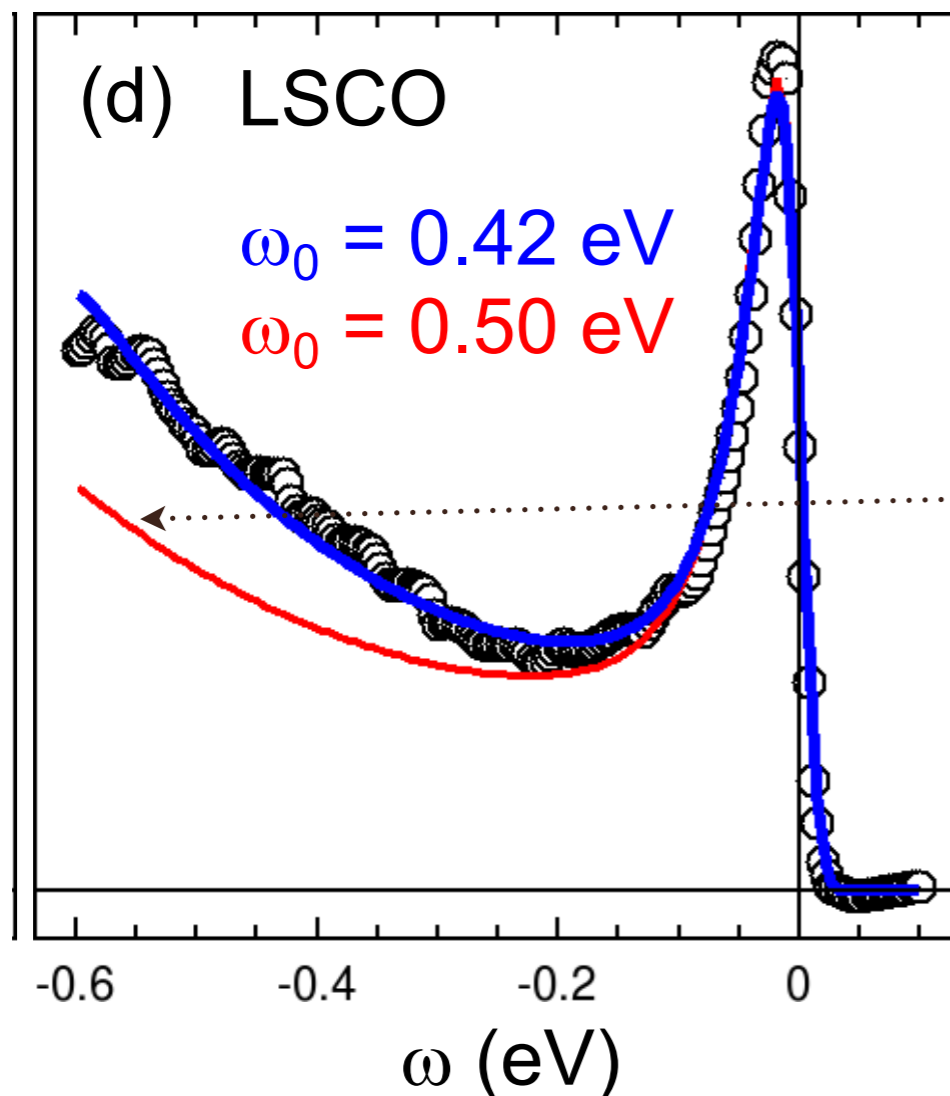
ECFL== Extremely Correlated Fermi liquid solution of tJ model.

$$G(k, z) = \frac{a_G}{z - \hat{E}_k - \Sigma(k, z)}, \quad \text{Dyson}$$

$$= \frac{a_G + \Psi(k, z)}{z - \hat{E}_k - \Phi(k, z)}. \quad \text{ECFL}$$

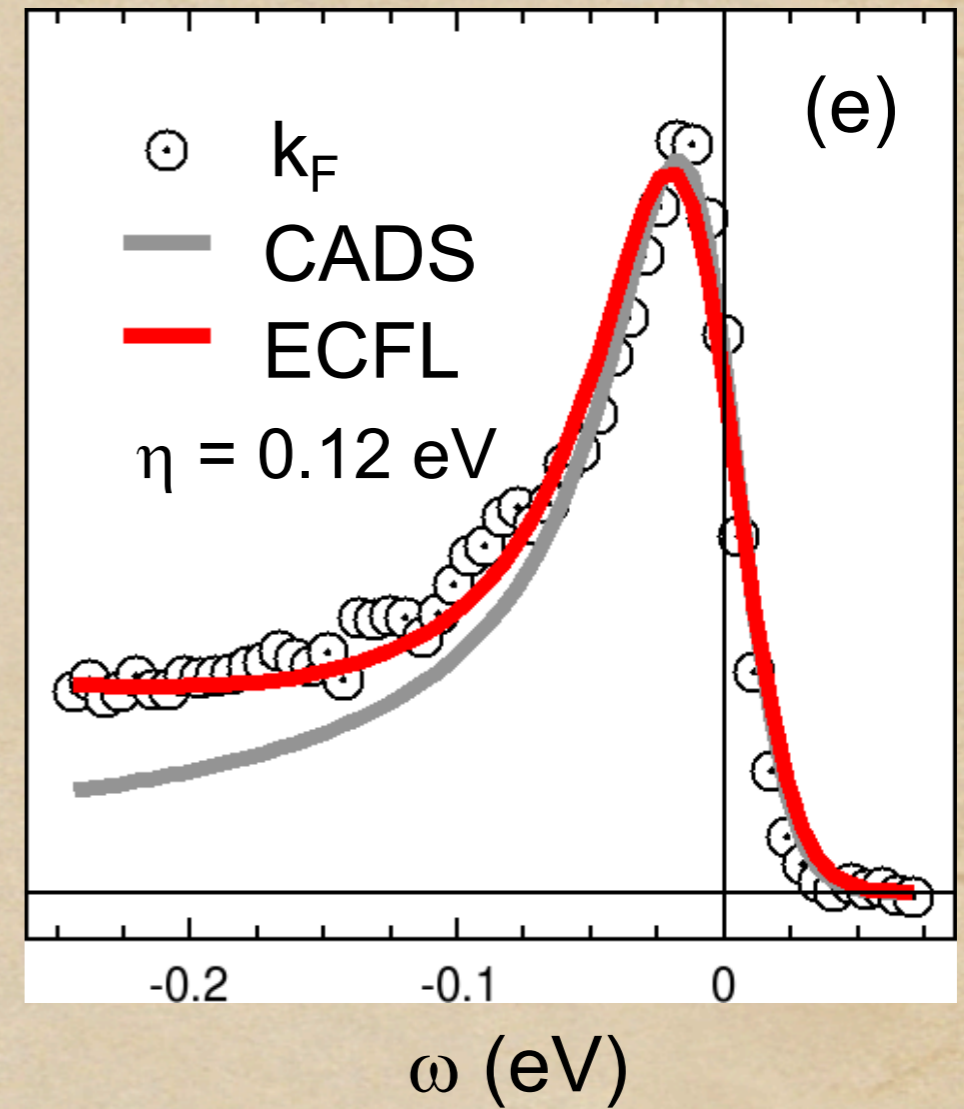
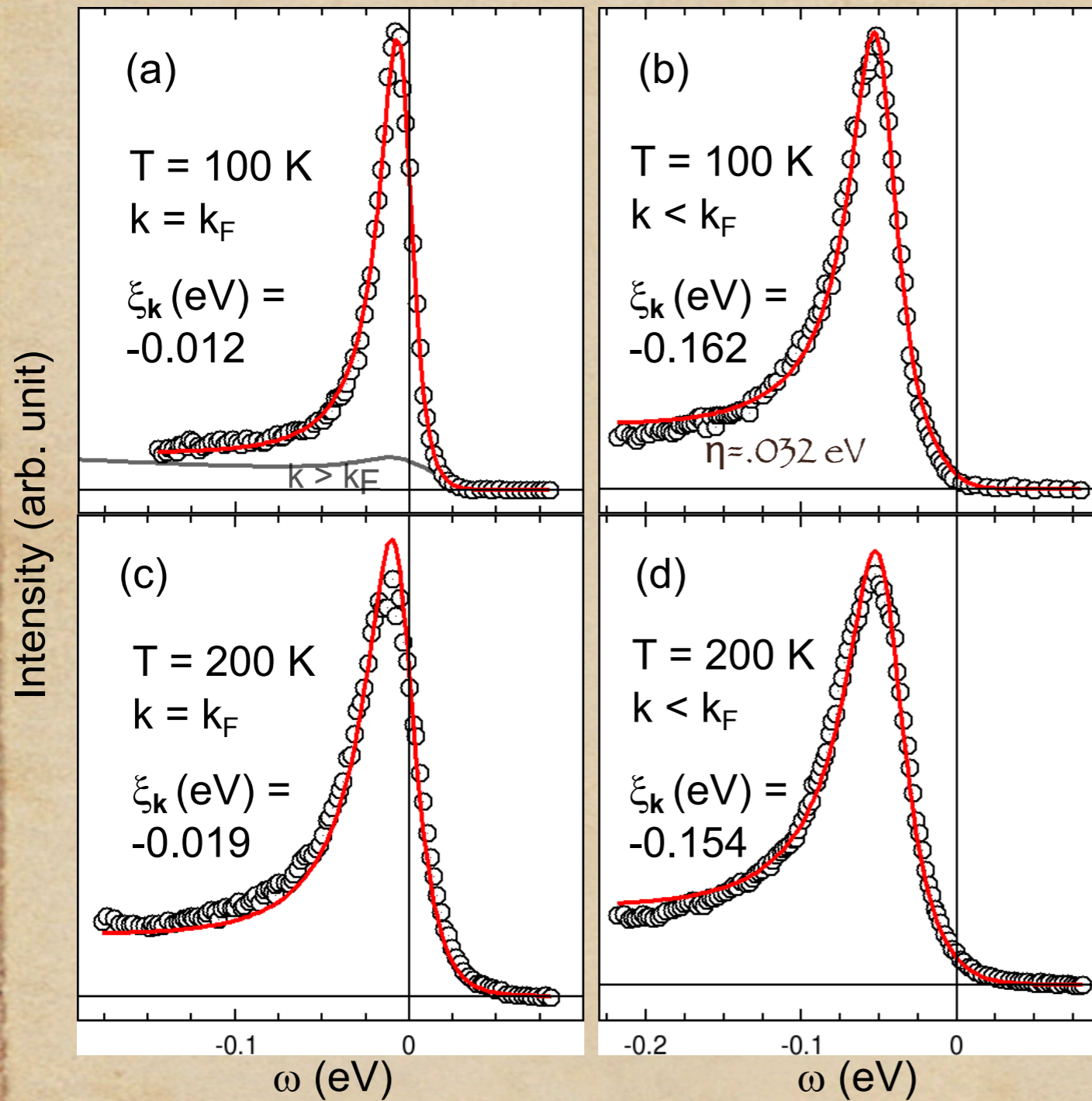
$$a_G = 1 - n/2, \quad n = \text{density}$$

extreme correlation model is best described in terms of "two self energies".



Smoking gun
 Linear rise of intensity
 for occupied states.

$$A(\vec{k}, \omega) = A_{FL}(\vec{k}, \omega) \left(1 - \frac{n}{2} + \frac{n^2}{4} \cdot \frac{\xi_{\vec{k}} - \omega}{\Delta_0} \right)_+$$



Same physical parameters only η different

η arises from impurities and sample penetration of ARPES.