

HW #3

Problem 1

$$\textcircled{a} \quad \frac{d\chi}{d(\ln \mu)} = b_0$$

$$\Rightarrow 8\pi^2 \frac{d}{d(\ln \mu)} \left(\frac{1}{g^2} \right) = b_0$$

$$\Rightarrow - \frac{16\pi^2}{g^3} \frac{dg}{d(\ln \mu)} = b_0$$

$$\Rightarrow \boxed{\frac{dg}{d \ln(\mu)} = - \frac{b_0 g^3}{16\pi^2}}$$

$$\textcircled{b} \quad \frac{d\chi}{d(\ln \mu)} = b_0 \Rightarrow \chi(\mu) - \chi(m) = b_0 [\ln(\mu) - \ln(m)]$$

$$\Leftrightarrow \boxed{\frac{1}{\alpha(\mu)} = \frac{1}{\alpha(m)} + \frac{b_0}{2\pi} \ln\left(\frac{\mu}{m}\right)}$$

$$\textcircled{c} \quad \alpha_1(\mu) = \alpha_2(m) \Rightarrow$$

$$\frac{1}{\alpha_1(m_2)} + \frac{b_0}{2\pi} \ln\left(\frac{\mu}{m_2}\right) = \frac{1}{\alpha_2(m_2)} + \frac{b_0}{2\pi} \ln\left(\frac{\mu}{m_2}\right)$$

$$\ln\left(\frac{\mu}{m_2}\right) = \frac{2\pi}{b_0^2 - b_0^2} \left\{ \frac{1}{\alpha_1(m_2)} - \frac{1}{\alpha_2(m_2)} \right\}$$

$$= \frac{2\pi \times 30.5}{7.016} \Rightarrow \boxed{\mu \approx 4 \times 10^{13} \text{ GeV}}$$

HW#3

$$\alpha_3(\mu) \approx 0.029$$

$$\& \alpha_1(\mu) = \alpha_2(\mu) \approx 0.023$$

$$\frac{1}{\alpha_1(\mu)} = \frac{1}{\alpha_2(\mu)} \approx \frac{1}{43.4}$$

$$\frac{1}{\alpha_3(\mu)} \approx 34.2$$

Long time ago I plotted

these 3 lines. \rightarrow It's attached at the end.

Problem 3.4

When pion π^- at rest $L=0$, $J=\underline{\underline{1}}$

thus

$$\text{L.H.S.} \quad (-1)^L P_{\pi^-} P_d = (-1)^0 (-1)(+1) = -1 =$$

R.H.S.

$$(-1)^L P_n P_n P_{\pi^0}$$

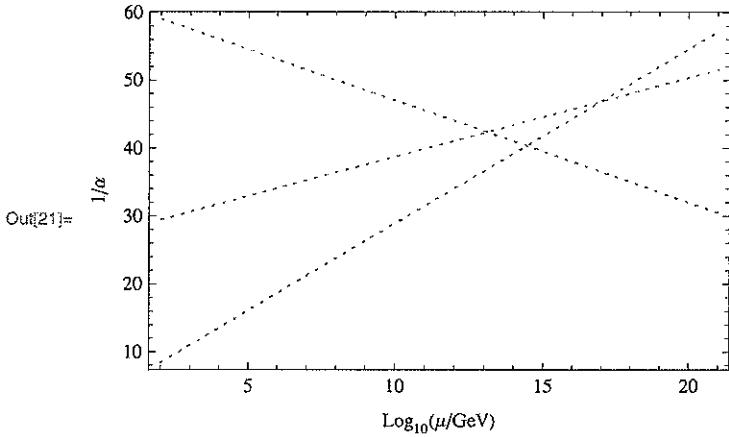
Now since neutrons are fermions the wave functions must be antisymmetric under neutron exchange. thus $L=1$

therefore

$$\text{R.H.S.} = (-1)^1 (+1)(+1)(-1) = +1 \quad \underline{\underline{1}}$$

Parity violation

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In[21]:= (*Standard Model *)gSM = Plot[{alphaInv1[x, alpha3Mz, b3sm],
  alphaInv1[x, alpha2Mz, b2sm], alphaInv1[x, 5 / 3 * alpha1Mz, b1sm]},
 {x, 2, 21}, Axes → False, Frame → True, FrameLabel → {"Log10(μ/GeV)", "1/α"}, PlotStyle → {{Gray, Dotted}, {Gray, Dotted}, {Gray, Dotted}}]
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```
In[22]:= stitchedAlpha1[alphaMz_, b_List, thr_List] :=
Module[{sm, alphaMSSM, mssm, xThr},
xThr = Log[10, thr];
sm = alphaInv1[x, alphaMz, b[[1]], thr[[1]]];
alphaMSSM = 1 / alphaInv1[xThr[[2]], alphaMz, b[[1]], thr[[1]]];
mssm = alphaInv1[x,
  alphaMSSM, b[[2]], thr[[2]]];
Piecewise[{{
sm, x < xThr[[2]]},
{mssm, xThr[[2]] ≤ x}
}]]
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In[23]:= stitchedAlpha2[alphaMz_, b_List, thr_List] :=
Module[{sm, alphaMSSM, mssm, alphaMSSM4hd, mssm4hd, xThr},
xThr = Log[10, thr];
sm = alphaInv1[x, alphaMz, b[[1]], thr[[1]]];
alphaMSSM = 1 / alphaInv1[xThr[[2]], alphaMz, b[[1]], thr[[1]]];
mssm = alphaInv1[x,
  alphaMSSM, b[[2]], thr[[2]]];
alphaMSSM4hd = 1 / alphaInv1[xThr[[3]],
  alphaMSSM, b[[2]], thr[[2]]];
mssm4hd = alphaInv1[x,
  alphaMSSM4hd, b[[3]], thr[[3]]];
Piecewise[{{
sm, x < xThr[[2]]},
{mssm, xThr[[2]] ≤ x ≤ xThr[[3]]},
{mssm4hd, xThr[[3]] ≤ x}
}]]
```

HW # 3

Perkins 3.9

Back of the book!

Problem 6

(a) Quarks have spin $\frac{1}{2}$ thus.

Meson spins are $= \frac{1}{2} + \frac{1}{2} = 1 \leftarrow$ vector meson

& $\frac{1}{2} - \frac{1}{2} = 0 \leftarrow$ scalar meson.

(b) Adding the first two quarks get me,

$$|\frac{1}{2}\rangle|\frac{1}{2}\rangle \rightarrow |1\rangle \& |0\rangle$$

& then, $|1\rangle|\frac{1}{2}\rangle \rightarrow |\frac{3}{2}\rangle \& |\frac{1}{2}\rangle||\frac{1}{2}\rangle$

also $|0\rangle|\frac{1}{2}\rangle \rightarrow |\frac{1}{2}\rangle \&$

\ Problem 7,

$$\frac{I_\mu(t)}{I_\mu(0)} = 1 - \sin^2 2\theta \sin^2 \left[\frac{(E_3 - E_2)t}{2} \right] \quad \dots \quad (1)$$

Now $E_i = p + \frac{m_i^2}{2p}$ from eqn. (4.9)

then, $E_3 - E_2 = \frac{1}{2p} [m_3^2 - m_2^2]$

Since $m^2 \ll E^2$ thus $E^2 \approx p^2 c^2$

then,

$$\frac{I_\mu(t)}{I_\mu(0)} \approx 1 - \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 t}{4E} \right)$$

Now $t = \frac{L}{c}$ also notice to fix the unit in
 (1) we need an \hbar in the denominator.

thus.

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta \cdot \sin^2 \theta \left(\frac{1}{4\hbar c} \frac{\Delta m^2 L}{E} \right)$$

Now $\hbar c = 0.197 \text{ GeV fm} = \cancel{0.197 \text{ MeV fm}}$
 ~~$= 0.197 \times 10^{+6} \text{ GeV fm}$~~

$$\hbar c = 0.197 \text{ GeV fm}$$

$$= 0.197 \times 10^{-6} \text{ eV m}$$

$$= 0.197 \times 10^{-9} \text{ eV km}$$

then. $4\hbar c = 0.788 \times 10^{-9} \text{ eV km}$

then. $\frac{1}{4\hbar c} = 1.27 \times 10^9 \text{ eV}^{-1} \text{ km}^{-1}$

$$\begin{aligned}\Delta m^2 &\rightarrow \text{eV}^2 \quad \& E \rightarrow \text{GeV} \\ &\rightarrow \text{eV} (10^9 \text{ GeV})\end{aligned}$$

Therefore -

we have.

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 \theta \sin^2 \theta \underbrace{\left(1 - 27 \frac{(\Delta m/\text{eV})^2 (\text{L}/\text{km})}{(E/\text{GeV})} \right)}_{\equiv}$$

Perkin 4.1

$$\Phi_u = 250 \text{ m}^2 \text{s}^{-1}$$

Rs

$$\text{Rate of ionization } R = 2.5 \text{ MeV} \frac{\text{cm}^2}{\text{gm}}$$

$$= 0.25 \text{ MeV kg}^{-1} \text{ m}^{-2}$$

then

$$\text{total dose} = \Phi_u \times R \times \text{Year.} \quad \text{approximate gray.}$$

$$= 3.2 \times 10^{-4} \text{ gray}$$

$$= 0.032 \text{ rad.}$$

One tenth of the natural dose.

$$\text{Now } 1 \text{ MeV} \approx 2 \times 10^{-30} \text{ kg}$$

so we have about 60×10^{30} protons.

Let the surface area $\approx 2 \text{ m}^2$

then # protons

$$\text{then } R_p = \frac{8000}{2 \times 10^{30}} \text{ MeV m}^{-2}$$

thus there's about 10^{30} MeV/kg radiation from proton decay. Note

$$\text{Then the } \text{mean} \text{ proton lifetime} = \frac{10^{30} \text{ MeV/kg}}{300 \text{ rad}} \times 80 \text{ yrs.}$$

> age of the universe.

You can do even better by considering the universe. \approx