Physics 220- Fall 2011 Theory of Many Body Physics Solution to Examination 2 Attempt any 4 questions 100 points, Time 1.5 hours 26 October, 2011

1. Electrons with spin half in 1 dimension live on a ring containing N sites. Find the relation between the number density of electrons $n = N_e/N$ and the Fermi momentum. What is the Fermi wave number at a density $n = 1/3$? [25]

We can write the number of particles including a factor of two for spin as

$$
n = \frac{N_e}{N} = \frac{2}{N} \sum_{k} \theta(k_F - |k|) \to 2 \int_{-k_F}^{k_F} \frac{dk}{2\pi} = 2\frac{k_F}{\pi},
$$

hence at $n = 1/3$ we find $k_F = \pi/6$.

2. The above electrons have a hopping matrix element t_0 at nearest neighbor separation, and t_1 at second neighbor separation. Find t_1/t_0 such that the net dispersion has a flat bottom near $k \sim 0$ going as $\varepsilon_k =$ $c_1 + c_2 k^4$ (i.e. missing the quadratic piece). [25] We can write the electronic dispersion as

$$
\varepsilon_k = -t_0 \left[e^{ik} + e^{-ik} \right] - t_1 \left[e^{i2k} + e^{-i2k} \right] = -2t_0 \cos(k) - 2t_1 \cos(2k),
$$

so expanding near $k \sim 0$ we get

$$
\varepsilon_k = -2t_0 - 2t_1 + (t_0 + 4t_1)k^2 - \frac{1}{12}(t_0 + 16t_1)k^4 + O(k^6).
$$

Setting $t_1 = -1/4t_0$ we get the required result

$$
\varepsilon_k = -(3/2)t_0 + \frac{1}{4}t_0 k^4 + O(k^6).
$$

3. Starting from the Fermionic version of the Jordan Wigner string

$$
J[1, n] = e^{i\pi \sum_{j=1}^{n-1} C_j^{\dagger} C_j},
$$

show that

$$
J[1, n]J[1, n] = 1.
$$

[25]

Since the eigenvalues of $C_j^{\dagger}C_j$ are either 1 or 0, we see that the value of the operator

$$
e^{i2\pi \sum_{j=1}^{n-1}C_j^{\dagger}C_j} \to 1.
$$

4. With two Fermionic destruction operators C_j and $j = 1, 2$ with the canonical properties $\{C_i, C_j\} = 0$ and $\{C_i, C_j^{\dagger}\} = \delta_{ij}$, we construct two other Fermi operators:

$$
\alpha = u C_1 + v C_2,
$$

$$
\beta = p C_1 + q C_2.
$$

Write down the conditions on u, v, p, q in order that the new Fermions are canonical and mutually independent. [25]

Taking u, v as real we need to satisfy the conditions

$$
\{\alpha,\alpha^{\dagger}\}=1=\{\beta,\beta^{\dagger}\}; \ \{\alpha,\beta^{\dagger}\}=0.
$$

These lead to

$$
u^{2} + v^{2} = 1 = p^{2} + q^{2},
$$

$$
u p + v q = 0.
$$

The usual choice that achieves these is to write

$$
u = \cos(\theta), v = \sin(\theta), p = -v, q = u.
$$

5. For spin half electrons on a lattice, the spin density is given as $S_j^z = \frac{1}{n} (n_x - n_y)$ and the number density as $a = (n_x + n_y)$. Show that the $\frac{1}{2}(n_{i\uparrow} - n_{i\downarrow})$, and the number density as $\rho_j = (n_{i\uparrow} + n_{i\downarrow})$. Show that the Hubbard interaction term

$$
U n_{i\uparrow} n_{i\downarrow},
$$

is expressible as a suitable linear combination of $U(S_i^z)^2$, and the number density $U\rho_i$.[25]

Let us start with

$$
(S_i^Z)^2 = \frac{1}{4}(n_{i\uparrow} - n_{i\downarrow})^2 = \frac{1}{4}(n_{i\uparrow}^2 + n_{i\downarrow}^2 - 2n_{i\uparrow}n_{i\downarrow}).
$$

On using

$$
n_{i\sigma}^2 = n_{i\sigma},
$$

 \boldsymbol{we} \boldsymbol{get}

$$
(S_i^Z)^2 = \frac{1}{4}(\rho_i - 2n_{i\uparrow}n_{i\downarrow}).
$$

 $Hence$

$$
U n_{i\uparrow} n_{i\downarrow} = -2U(S_i^Z)^2 + \frac{U}{2}\rho_i.
$$