Physics 220- Fall 2011 Theory of Many Body Physics Solution to Examination 2 Attempt any 4 questions 100 points, Time 1.5 hours 26 October, 2011

1. Electrons with spin half in 1 dimension live on a ring containing N sites. Find the relation between the number density of electrons $n = N_e/N$ and the Fermi momentum. What is the Fermi wave number at a density n = 1/3? [25]

We can write the number of particles including a factor of two for spin as

$$n = \frac{N_e}{N} = \frac{2}{N} \sum_{k} \theta(k_F - |k|) \to 2 \int_{-k_F}^{k_F} \frac{dk}{2\pi} = 2\frac{k_F}{\pi},$$

hence at n = 1/3 we find $k_F = \pi/6$.

2. The above electrons have a hopping matrix element t_0 at nearest neighbor separation, and t_1 at second neighbor separation. Find t_1/t_0 such that the net dispersion has a flat bottom near $k \sim 0$ going as $\varepsilon_k = c_1 + c_2 k^4$ (i.e. missing the quadratic piece). [25] We can write the electronic dispersion as

$$\varepsilon_k = -t_0 \left[e^{ik} + e^{-ik} \right] - t_1 \left[e^{i2k} + e^{-i2k} \right] = -2t_0 \cos(k) - 2t_1 \cos(2k),$$

so expanding near $k \sim 0$ we get

$$\varepsilon_k = -2t_0 - 2t_1 + (t_0 + 4t_1)k^2 - \frac{1}{12}(t_0 + 16t_1)k^4 + O(k^6).$$

Setting $t_1 = -1/4t_0$ we get the required result

$$\varepsilon_k = -(3/2)t_0 + \frac{1}{4}t_0k^4 + O(k^6).$$

3. Starting from the Fermionic version of the Jordan Wigner string

$$J[1,n] = e^{i\pi \sum_{j=1}^{n-1} C_j^{\dagger} C_j},$$

show that

$$J[1,n]J[1,n] = 1.$$

[25]

Since the eigenvalues of $C_j^{\dagger}C_j$ are either 1 or 0, we see that the value of the operator

$$e^{i2\pi\sum_{j=1}^{n-1}C_j^{\dagger}C_j} \to 1.$$

4. With two Fermionic destruction operators C_j and j = 1, 2 with the canonical properties $\{C_i, C_j\} = 0$ and $\{C_i, C_j^{\dagger}\} = \delta_{ij}$, we construct two other Fermi operators:

$$\alpha = u C_1 + v C_2,$$

$$\beta = p C_1 + q C_2.$$

Write down the conditions on u, v, p, q in order that the new Fermions are canonical and mutually independent. [25]

Taking u, v as real we need to satisfy the conditions

$$\left\{\alpha, \alpha^{\dagger}\right\} = 1 = \left\{\beta, \beta^{\dagger}\right\}; \ \left\{\alpha, \beta^{\dagger}\right\} = 0.$$

These lead to

$$u^{2} + v^{2} = 1 = p^{2} + q^{2},$$

 $u \ p + v \ q = 0.$

The usual choice that achieves these is to write

$$u = \cos(\theta), v = \sin(\theta), p = -v, q = u.$$

5. For spin half electrons on a lattice, the spin density is given as $S_j^z = \frac{1}{2}(n_{i\uparrow} - n_{i\downarrow})$, and the number density as $\rho_j = (n_{i\uparrow} + n_{i\downarrow})$. Show that the Hubbard interaction term

$$U n_{i\uparrow} n_{i\downarrow},$$

is expressible as a suitable linear combination of $U(S_i^z)^2$, and the number density $U\rho_i$.[25]

Let us start with

$$(S_i^Z)^2 = \frac{1}{4}(n_{i\uparrow} - n_{i\downarrow})^2 = \frac{1}{4}(n_{i\uparrow}^2 + n_{i\downarrow}^2 - 2n_{i\uparrow}n_{i\downarrow}).$$

On using

$$n_{i\sigma}^2 = n_{i\sigma},$$

 $we \ get$

$$(S_i^Z)^2 = \frac{1}{4}(\rho_i - 2n_{i\uparrow}n_{i\downarrow}).$$

Hence

$$U n_{i\uparrow} n_{i\downarrow} = -2U(S_i^Z)^2 + \frac{U}{2}\rho_i.$$